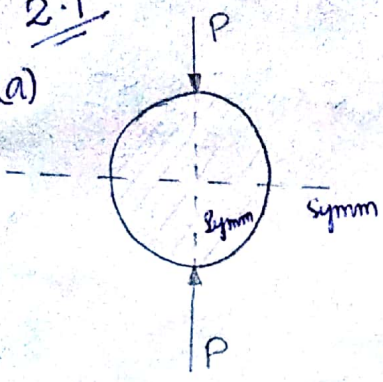
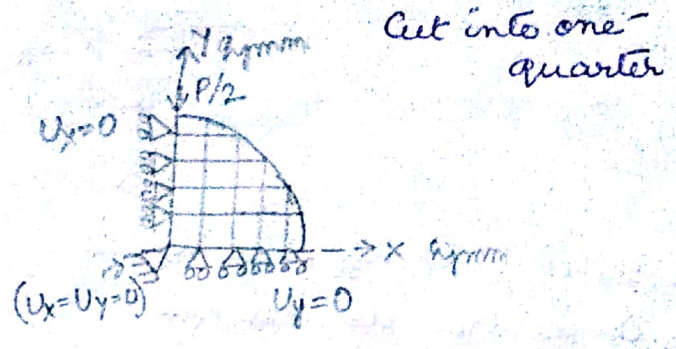


2:1

(a)

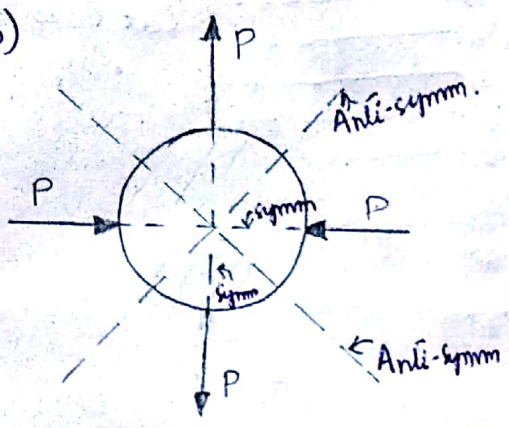


Circular disk with 2 diametrically opposite force

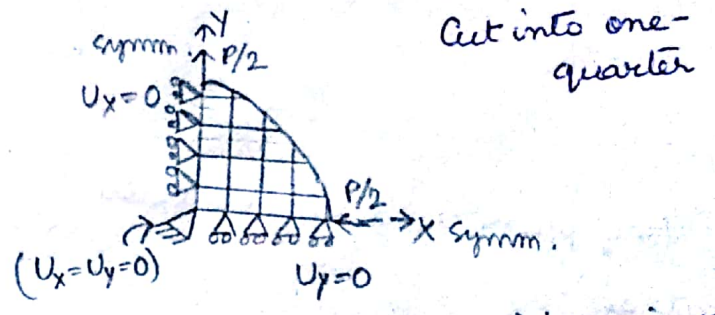


Cut into one-quarter

(b)



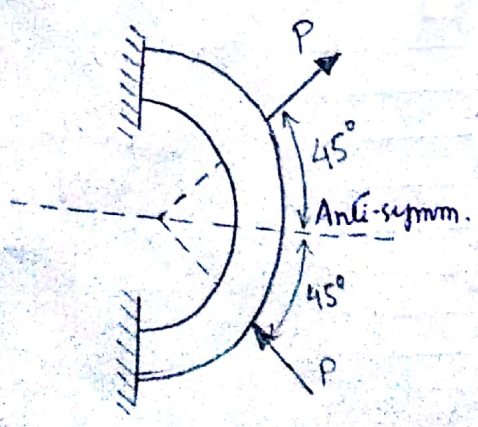
Circular disk with 2 diametrically opposite force pairs



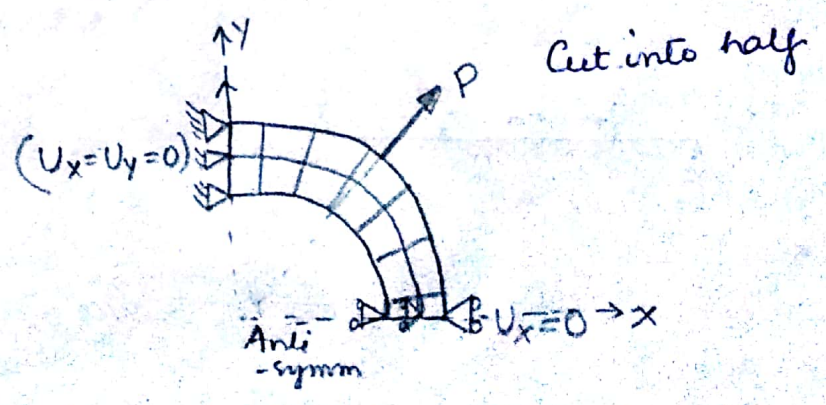
Cut into one-quarter

This can be further divided into a one-eighth part with zero displacement along anti-symmetric line.

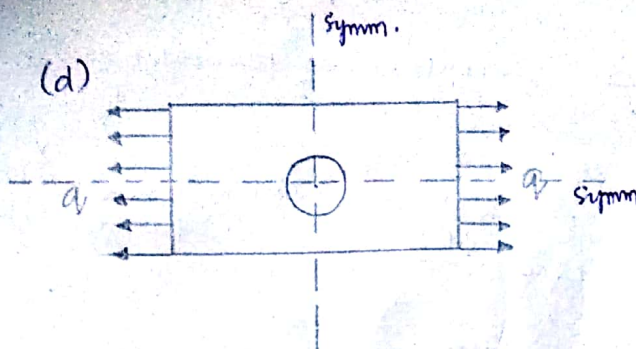
(c)



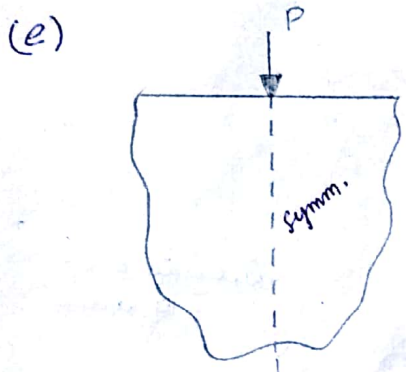
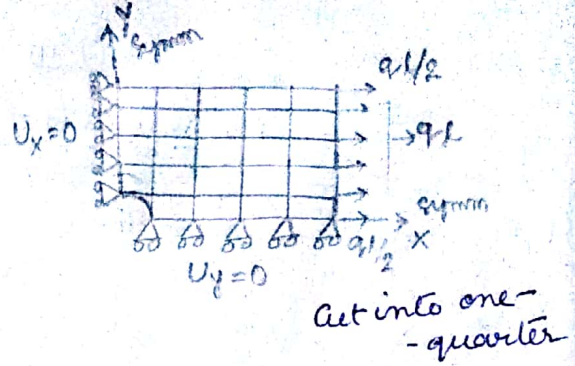
Clamped semiannulus



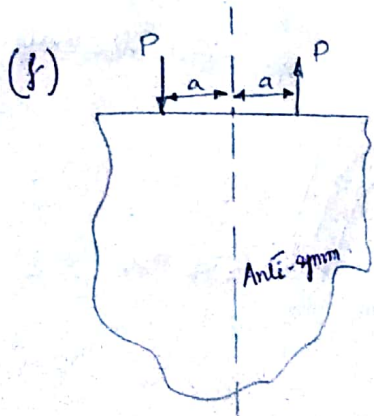
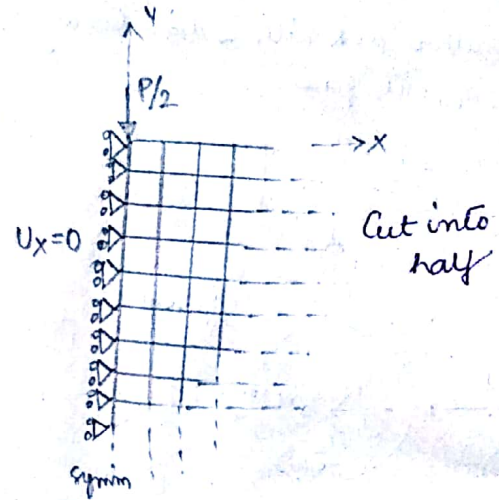
Cut into half



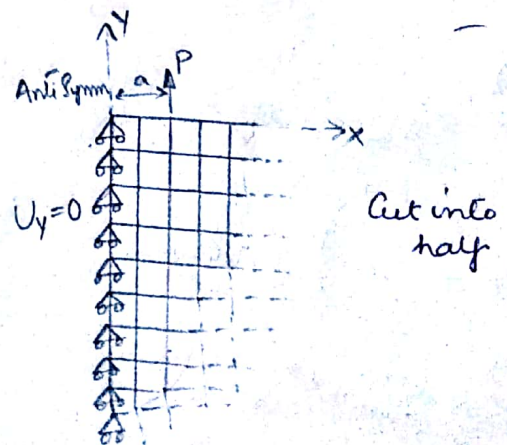
Stretched rectangular plate with central circular hole



Half plane under concentrated load

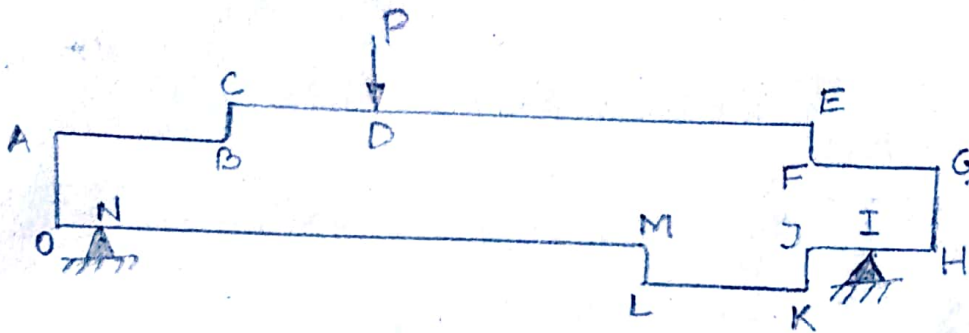


Half plane under concentrated load



* We need to put a constraint in the infinite plates in (e) & (f) as finite element analysis favors finite length models & not infinite models.

2.2



The spots in the above figure that would require a locally finer mesh are as follows -

- (i) Point - D (Concentrated load at D)
- (ii) Points - N & I (Points where reaction forces will act)
- (iii) Points - B, F, J & M (Re-entrant points - although they are not sharp edged, still they are likely to have a high stress concentration)

2.3

$$A = A_i(1 - \xi_3) + A_j \xi_3$$

$$q(x) = \rho A \omega^2 x$$

$$\therefore \text{Nodal force, } f^e = \rho \int_0^L \begin{bmatrix} 1 - \xi_3 \\ \xi_3 \end{bmatrix} dx$$

$$= \int_0^L \rho A \omega^2 x \begin{bmatrix} 1 - \xi_3 \\ \xi_3 \end{bmatrix} dx \quad \text{--- (1)}$$

Now,

$$\xi_3 = \frac{x - x_i}{L}$$

$$\Rightarrow x = \xi_3 L + x_i$$

$$\Rightarrow dx = L d\xi_3$$

\therefore From (1),

$$f^e = \int_0^1 \rho A \omega^2 (\xi_3 L + x_i) \begin{bmatrix} 1 - \xi_3 \\ \xi_3 \end{bmatrix} L d\xi_3$$

For $x_i = 0$,

$$f^e = \int_0^1 \rho A \omega^2 (\xi_3 L) \begin{bmatrix} 1 - \xi_3 \\ \xi_3 \end{bmatrix} L d\xi_3$$

Substituting the value of A, we get

$$f^e = \int_0^1 \rho \omega^2 L^2 [A_i(1 - \xi_3) + A_j \xi_3] \xi_3 \begin{bmatrix} 1 - \xi_3 \\ \xi_3 \end{bmatrix} d\xi_3$$

$$= \int_0^1 \rho \omega^2 L^2 [A_i(1 - \xi_3) + A_j \xi_3] \begin{bmatrix} \xi_3 - \xi_3^2 \\ \xi_3^2 \end{bmatrix} d\xi_3$$

$$= \int_0^1 \rho \omega^2 L^2 \begin{bmatrix} A_i \xi_3 - \xi_3^2 (2A_i - A_j) + \xi_3^3 (A_i - A_j) \\ A_i (\xi_3^2 - \xi_3^3) + A_j \xi_3^3 \end{bmatrix} d\xi_3$$

$$= \rho \omega^2 L^2 \begin{bmatrix} A_i \frac{\xi_3^2}{2} - \frac{\xi_3^3}{3} (2A_i - A_j) + \frac{\xi_3^4}{4} (A_i - A_j) \\ A_i (\xi_3^3/3 - \xi_3^4/4) + A_j \xi_3^4/4 \end{bmatrix}_0^1$$

$$= \rho \omega^2 l^2 \begin{bmatrix} \frac{A_i}{2} - \left(\frac{2A_i - A_j}{3} \right) + \left(\frac{A_i - A_j}{4} \right) \\ A_i \left(\frac{1}{3} - \frac{1}{4} \right) + A_j/4 \end{bmatrix}$$

$$f^e = \rho \omega^2 l^2 \begin{bmatrix} \frac{1}{12} (A_i + A_j) \\ \frac{A_i}{12} + A_j/4 \end{bmatrix}$$

This is the required expression of consistent nodal forces.

For a prismatic bar,

$$A = A_i = A_j$$

$$\therefore f^e = \rho \omega^2 l^2 \begin{bmatrix} \frac{1}{12} (A + A) \\ A \left(\frac{1}{12} + \frac{1}{4} \right) \end{bmatrix}$$

$$f^e = \rho \omega^2 l^2 \begin{bmatrix} A/6 \\ A/3 \end{bmatrix}$$