

FEF Exam

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Mid - Term

Computational Mechanics

Ex-1

$$u_t + u u_x = 0 \quad \text{--- (1)}$$

$$u = u_0(x)$$

$$u = 1$$

$$u_0(x) = \begin{cases} 1 & , x \in [0, p] \\ 1 - \frac{x-p}{q-p} & , x \in [p, q] \\ 0 & , x \in [q, 1] \end{cases}$$

$$p = 0.64 \quad \text{and} \quad q = 0.84$$

$$f_x(u) = u u_x \quad \text{--- (2)}$$

$$u_t + f_x(u) = 0 \quad \text{(Cauchy eq)} \quad \text{--- (3)}$$

$$u = u_0(x)$$

Taylor series

$$u^{n+1} = u^n + \Delta t u_t^n + \frac{(\Delta t)^2}{2} u_{tt}^n + \underbrace{O(\Delta t)^3}_{\text{truncation error}}$$

$$\frac{u^{n+1} - u^n}{\Delta t} = u_t^n + \frac{\Delta t}{2} u_{tt}^n + \underbrace{O(\Delta t)^2}_{\text{Neglect this error}}$$

We get

$$\frac{u^{n+1} - u^n}{\Delta t} = u_t^n + \frac{\Delta t}{2} u_{tt}^n \quad \text{--- (4)}$$

$$u_t = -f_x \quad \text{(From Cauchy)} \quad \text{--- (5)}$$

$$u_{tt} = -f_{xt} = -f_{tx} = -(f u u_x)_x$$

Burger Eq. Time discretisation

$$\frac{u^{n+1} - u^n}{\Delta t} = -f_x^n + \frac{\Delta t}{2} ((u^n)^2 u_x^n)_x$$

Now, integration by parts of spatial terms
Tri-2 formulation is.

$$\int_0^L w \left(\frac{u^{n+1} - u^n}{\Delta t} \right) dx = \int_0^L w_x f^n dx - \frac{\Delta t}{2} \int_0^L w_x (u^n)^2 u_x^n dx$$

$$= \int_0^L w f_x^n dx - \frac{\Delta t}{2} \int_0^L w_x (u^n)^2 u_x^n dx$$

$$f_x = \frac{df}{du} \frac{\partial u}{\partial x} = -u^2 u_x$$

∴ ~~gth term~~ term of R.H.S, i.e.,

$$- \left[-w \frac{\Delta t}{2} (u^n)^2 u_x^n \right]_{x=0}^{x=L} = - \left[w \frac{\Delta t}{2} f_x^n \right]_0^L$$

Now Inlet Boundary Condⁿ is

$$f(0, t) = \left(\frac{u^2}{2} \right)_{x=0} = \frac{1}{2} [u(0, t) = 1]$$

3. Classical representation (2-pt gaussian quadrature)

~~Σ Ni~~

$$f = \sum f(N_i u_i)$$

Group Representation (Interpolation at ^{element}_n nodes)

$$f = \sum N_i f(u_i)$$

Ex 2 $-\gamma u_{xx} + \beta u_x = 0, x \in (-1, 1)$ — ①

$u = 0, x = -1$

$u = -1, x = 1$

$\gamma = 0.03$

$\beta = 1.8$ (convection vel.)

Solⁿ

Here $s = 0$ (source term)

Integrating eq ① and multiply by weights, w

$\int_{\Omega} w \beta u_x + \int_{\Omega} w \gamma u_{xx} = 0$ — ②

I II

$\int w \gamma u_{xx} = - \int w \beta u_x$

II term can be written as

$-\int w \nabla \cdot (\gamma \nabla u) =$

using divergence theorem

$\nabla \cdot (w \cdot (\gamma \nabla u)) = \nabla w \cdot (\gamma \nabla u) + \underbrace{w \nabla \cdot (\gamma \nabla u)}$

$\Rightarrow - \int w \nabla \cdot (\gamma \nabla u) = \int \nabla w \cdot (\gamma \nabla u) - \int \nabla \cdot (w \cdot (\gamma \nabla u))$

$w = 0$ for Γ_0

$\Rightarrow - \int w \nabla \cdot (\gamma \nabla u) = \int \nabla w \cdot \gamma \nabla u$

$= \int w_x \gamma u_x$ — ② ③

∴ Eq ② becomes.

$\int_{x=-1}^{x=1} w \beta u_x dx + \int_{x=-1}^{x=1} w_x \gamma u_x dx = 0$

$a(w, u) + c(w, \beta, u) = 0$

$\frac{\partial}{\partial x} (\gamma u_x) + \beta u_x$

$$\int_0^L \sum_{B=2}^{n_{eq}+1} \left(\beta N_A \frac{\partial N_B}{\partial x} + \gamma \frac{\partial N_A}{\partial x} \frac{\partial N_B}{\partial x} \right) u_B dx = 0$$

$$\text{Nodes} = \{1, n_{nd}\}$$

$$= \{1, 2, 3, \dots, n_{eq}, n_{eq}+1, n_{np}\}$$

$$N_1 = \frac{1}{2}(1-\xi); \quad N_2 = \frac{1}{2}(1+\xi); \quad dx = \frac{(x_2-x_1)}{2} d\xi = \frac{h}{2} d\xi$$

Element convection matrix

$$C^e = \beta \int_{\Omega^{(e)}} \begin{pmatrix} N_1 \frac{\partial N_1}{\partial x} & N_1 \frac{\partial N_2}{\partial x} \\ N_2 \frac{\partial N_1}{\partial x} & N_2 \frac{\partial N_2}{\partial x} \end{pmatrix} dx = \frac{\beta}{2} \begin{pmatrix} -1 & +1 \\ 1 & +1 \end{pmatrix}$$

Element diffusion matrix

$$K^{(e)} = \gamma \int_{\Omega^{(e)}} \begin{pmatrix} \frac{\partial N_1}{\partial x} \frac{\partial N_1}{\partial x} & \frac{\partial N_1}{\partial x} \frac{\partial N_2}{\partial x} \\ \frac{\partial N_2}{\partial x} \frac{\partial N_1}{\partial x} & \frac{\partial N_2}{\partial x} \frac{\partial N_2}{\partial x} \end{pmatrix} dx$$

$$= \frac{\gamma}{h} \begin{pmatrix} 1 & -1 \\ -1 & 1 \end{pmatrix}$$

b) There will be oscillations, in the eq or conv-diff eq. To reduce, we introduce stabilisation term

Stabilisation term

$$\sum_e \int P(w) \tau R(u)$$

$$P(w) = \beta \cdot \nabla w + \nabla \cdot (\gamma \nabla w) \\ = \beta w_x + \gamma w_{xx}$$

$$R(u) = \beta \cdot \nabla u - \nabla \cdot (\gamma \nabla u) \quad \xrightarrow{-s} 0 \\ = \beta u_x - \gamma u_{xx}$$

∴ stabilisation term is

$$\sum_e \int (\beta w_x + \gamma w_{xx}) \tau (\beta u_x - \gamma u_{xx})$$

Discretised form

$$a(w, u) + c(w, \beta, u) + \sum_e \int_{\Omega^e} (\beta w_x + \gamma w_{xx}) \tau (\beta u_x - \gamma u_{xx}) = 0$$

$$\int_{-1}^{+1} \gamma \frac{\partial N_A}{\partial x} \frac{\partial N_B}{\partial x} u_B + \int_{-1}^{+1} \beta N_A \frac{\partial N_B}{\partial x} u_B +$$

$$\sum_e \int_{\Omega^e} \left(\beta \frac{\partial N_A}{\partial x} + \gamma \frac{\partial^2 N_A}{\partial x^2} \right) \tau \left(\beta \frac{\partial N_B}{\partial x} - \gamma \frac{\partial N_B}{\partial x} \right) u_B = 0$$

SGS have no instabilities like SUPG and GLLS.

(c) Stabilization term: for linear element diffusion term is zero. For quadratic element $\gamma \neq 0$

For 1-D linear elements superconvergence is obtained.

τ vanishes when mesh is refined

Convergence depends on asymptotic behaviour of τ

However, 4th order accuracy is not reached for higher dimensions.

$$\tau = \frac{h}{2\beta} \left(\coth Pe - \frac{1}{Pe} \right)$$